

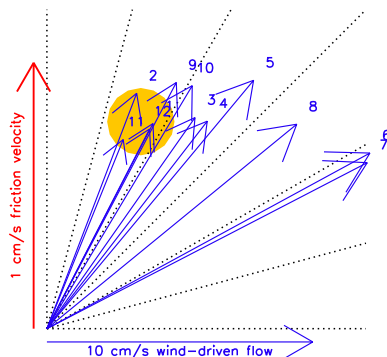
ストークスドリフトについて

吉川 裕

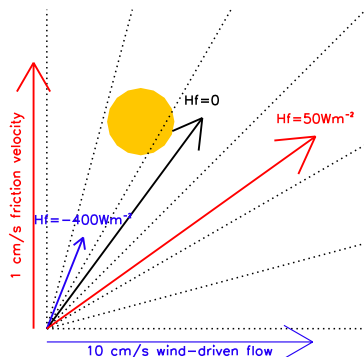
九大応力研

海洋レーダ研究集会

実測 (月平均値)



数値実験 (風+月平均熱フラックス)



- 実測と数値実験の差は熱フラックスの日変化が原因と考えている
- 波浪（ストークスドリフト）の影響は？と良く聞かれる

観測海域

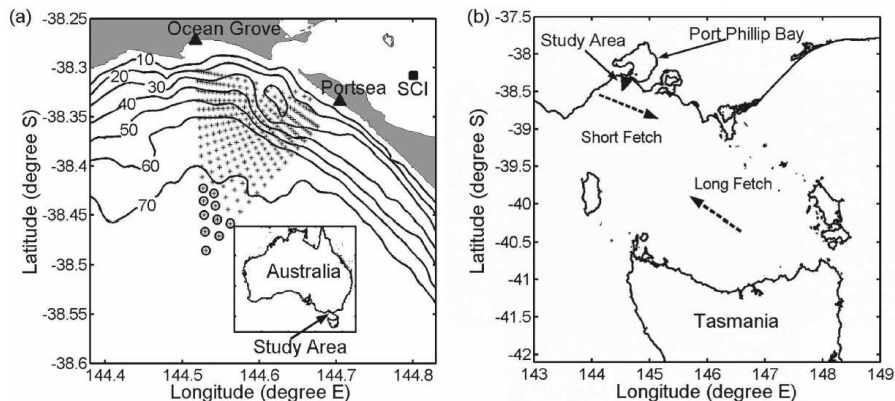


FIG. 1. Map of the study area (a) land (shaded), data grid points (asterisks), and bathymetry contours (m); two radar stations are marked as triangles, weather station SCI is marked as a square. The 10 grid points highlighted with circles are the ones chosen for the fetch study; (b) the study area in wider perspective shows how fetch varies with wind direction. The two dashed arrows represent the mean direction of wind during the short- and long-fetch duration in the analysis.

レーダ計測流と風

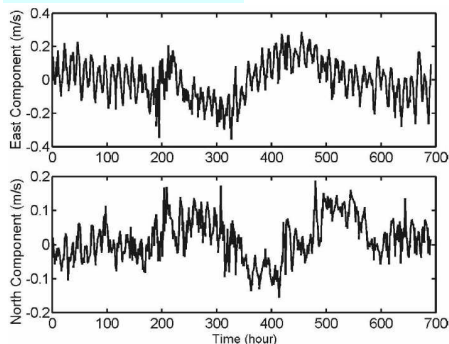


FIG. 2. Time series of HF radar current data at one of the southern grid points, reference time is 1700 Australian eastern standard time (UTC+10) 27 Jun 2001: (top) east and (bottom) north component.

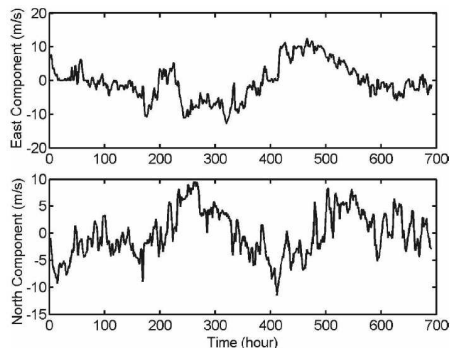


FIG. 3. Wind data from an anemometer located on SCI, reference time is 1700 Australian eastern standard time (UTC+10) 27 Jun 2001: (top) east and (bottom) north component.

レーダ計測流と風の相関係数と偏角

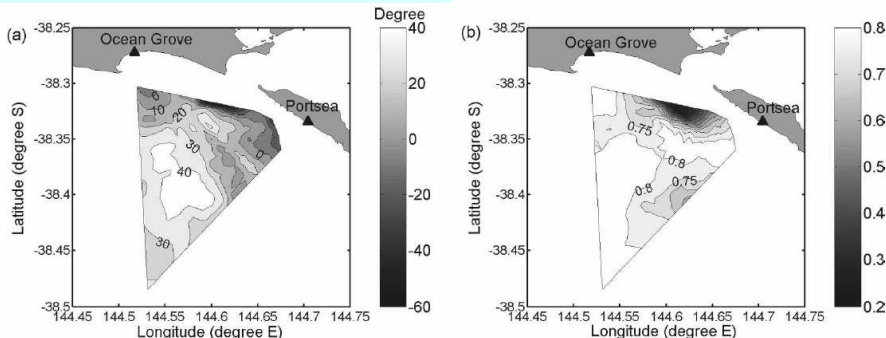


FIG. 6. Contours related to the response of surface current to wind over the whole time series; (a) the ϕ value is optimized for maximum correlation. Positive angle indicates that the responding current is on the left of the wind. (b) The correlation coefficient between the surface current and wind corresponding to the optimized direction relation in (a) is given.

解析期間

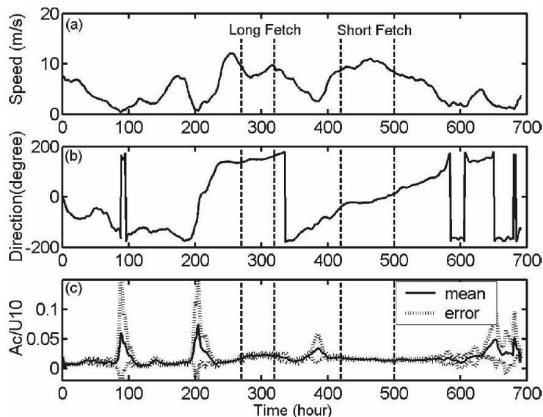


FIG. 7. The selected long- and short-fetch periods from (a) time series of the filtered wind speed, (b) time series of the angle of the filtered wind vector measured anticlockwise from east, and (c) time series of the average ratio of current speed to wind speed (solid line) with error bar (dashed line). The reference time is 1700 Australian eastern standard time (UTC + 10) 27 Jun 2001.

レーダ計測流と風 (Long Fetch と Short Fetch)

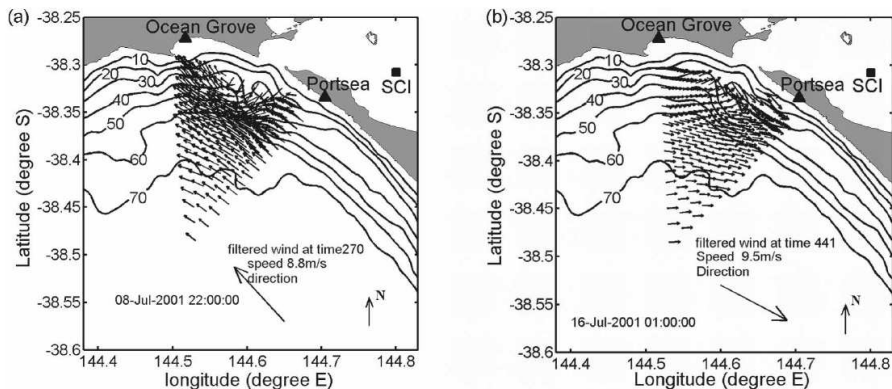


FIG. 8. Typical map of filtered surface currents under (a) long- and (b) short-fetch conditions.

結果

期間 (h)	Fetch	風力係数 (%)	偏角 (deg)
270-320	Long	2.1	14.4±3.8
420-500	Short	1.5	24.4±4.9

結果

期間 (h)	Fetch	風力係数 (%)	偏角 (deg)
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解釈

計測流はエクマン流とストークスドリフトの重ね合わせ

- エクマン流は風応力に比例し偏角は 45 度
- ストークスドリフトは風速に比例し偏角は 0

$$\underbrace{w (= u + iv)}_{\text{計測流}} = \underbrace{w_{EK}}_{\text{Ekman 流}} + \underbrace{w_{ST}}_{\text{Stokes 流}}$$

$$\propto U_{10}^2 \exp(i45^\circ) \propto U_{10}$$

$$\text{比例係数} \quad 6 \times 10^{-4} \quad 0.015 \quad (\text{Long Fetch})$$

$$6 \times 10^{-4} \quad 0.008 \quad (\text{Short Fetch})$$

レーダ計測流と推定値

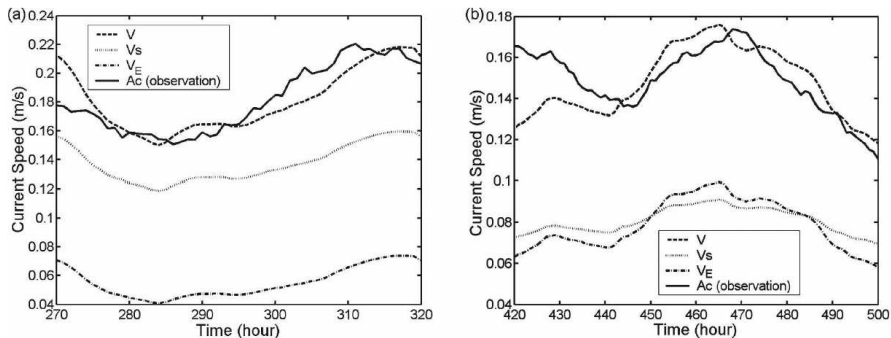


FIG. 12. Model results at grid point 255 for (a) long- and (b) short-fetch periods; the total surface current speed V (dashed), the observed surface current speed A_C (solid), speed of the Ekman-type current V_E (dash dotted), and Stokes drift V_S (dotted).

ストークスドリフトと Fetch

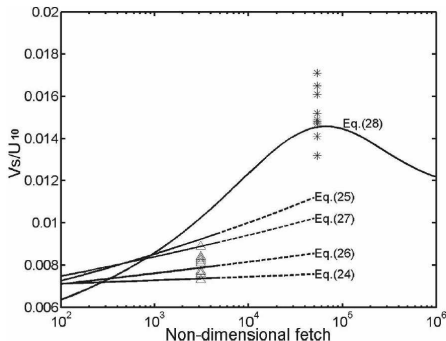


FIG. 13. The value of V_s/U_{10} from empirical functions and from our data; values from (24)–(27) are shown as combined solid and dashed lines (dashed lines represent regions where these equations do not apply). For the transition zone, the CERC (1973) line [Eq. (28)] is shown (solid). The asterisks (long fetch) and triangles (short fetch) represent the value from our data analysis (a_2 in Table 1).

- エクマン流 : 流速 $\propto$ 摩擦速度 ($\propto U_{10}$)、偏角 $\neq 45^\circ$
- 計測流 = エクマン流 + ストークスドリフト ?
- そもそも波浪の影響 = ストークスドリフト ?

回転系でのストークスドリフト

回転系・非粘性ではストークスドリフト = 0

(Ursell(1950)、Pollard(1970)、Hasselmann(1970)、...)

Jenkins (1986) (ラグランジュ座標系での解析解)

強粘性

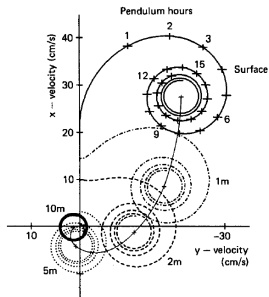


FIG. 4. Development in time (0 to 50 pendulum hours) of the mass transport velocity at different depths. No wind; wave height = 1.76 m for $t > 0$; eddy viscosity = $10^{-3} \text{ m}^2 \text{ s}^{-1}$.

弱粘性

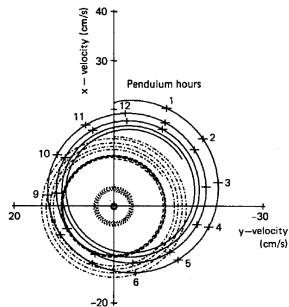


FIG. 5. Development in time (0 to 50 pendulum hours) of the mass transport velocity at different depths. Outermost to innermost: surface, 1, 2, 5 and 10 m. No wind; wave height = 1.76 m for $t > 0$; eddy viscosity = $10^{-5} \text{ m}^2 \text{ s}^{-1}$.

Xu and Bowen (1994) (オイラー座標系での解釈)

回転系での深水波 (正の x 方向に進行)

$$u \propto \frac{\omega}{\sqrt{1-\epsilon}} \omega A \cos(kx - \omega t) e^{kz}$$

$$w \propto \omega A \sin(kx - \omega t) e^{kz}$$

$$v \propto \frac{\epsilon}{\sqrt{1-\epsilon}} \omega A \sin(kx - \omega t) e^{kz} \quad \epsilon = \frac{f}{\omega} \simeq O(10^{-4})$$

波浪に伴う運動量輸送 (応力)

$$-\frac{\partial \overline{vw}}{\partial z} = -f \omega A^2 k e^{2kz} = -f u_{ST}^{IR,IV}$$

ストークスドリフトと大きさ同じで逆向きの平均流が波浪により生成

オイラー座標系での解釈: $u_{ST}^{R,V} = u_{ST}^{IR,IV} + u_{WM} = 0$

- 粘性力の重要性

⇒ $|u_{ST}^{R,V}| \neq |u_{ST}^{IR,IV}| = \omega A^2 k e^{2kz}$

- コリオリ力の重要性

⇒ $ANG(u_{ST}^{R,V}) \neq ANG(U_{10})$

- 波浪はラングミュア循環を引き起こすと言われる。

⇒ 鉛直混合過程も関与（もっと複雑）

対馬海峡での表層流に及ぼす波浪の影響（途中経過）

- 地衡流の直接評価 ⇨ 非地衡流の高精度推定
- 膨大な資料数 ⇨ 統計的有意性

海洋レーダー流速

- 除潮 x 成分流速を格子化
(0.05×0.05 deg)
- 空間平均 (灰色部)

水位 (JODC)

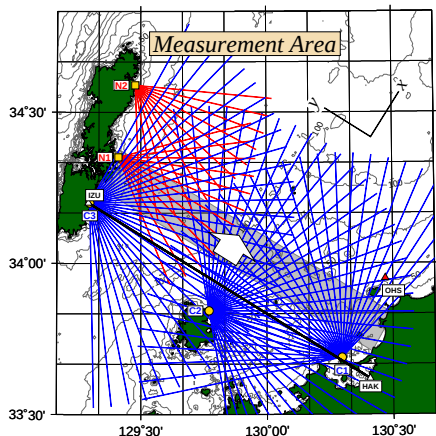
- 博多・厳原 (120 km)
- 毎時水位を除潮

摩擦速度 (JMA の解析値)

- 空間平均 (灰色部)
- 3 または 6 時間毎

波浪 (NOWPHAS)

- 大島
- 2 時間毎
- 有義波高、有義波周期、波向



Crossed Loop Antenna (CODAR)

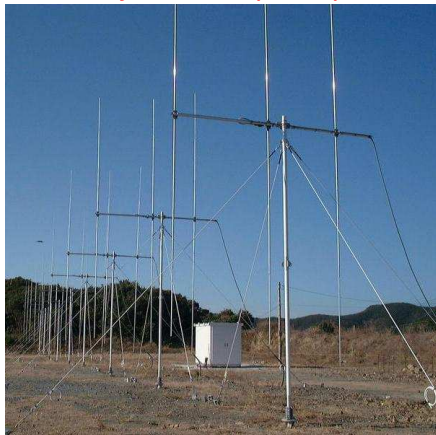


C5; OUMI Site (Tsushima Island)

CODAR (C1 ~ C5)

Center Frequency	13.9MHz
Range Coverage	75km
Range Resolution	3.0km
Bearing Coverage	360 deg
Bearing Resolution	5.0 deg
Sampling Time	4 min 16 s
Rad. Vel. Interval	1 hour
Measurement Depth	~ 172cm
Resonant Wave Length	11 m

DBF Array Antenna (NJRC)

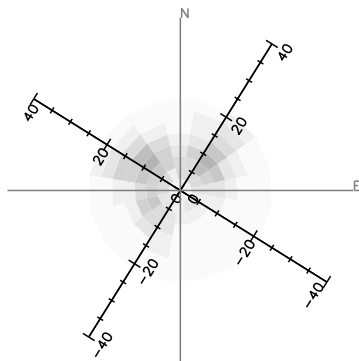


NJRC (N1,N2)

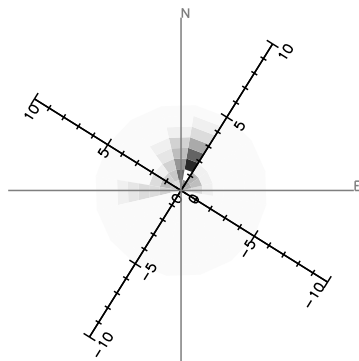
Center Frequency	24.5MHz
Range Coverage	50km
Range Resolution	1.5km
Bearing Coverage	90 deg
Bearing Resolution	7.5 deg
Sampling Time	30 min.
Rad. Vel. Interval	1 hour
Measurement Depth	~ 98cm
Resonant Wave Length	6.1 m

N1; Akashima Site (Tsushima Island)

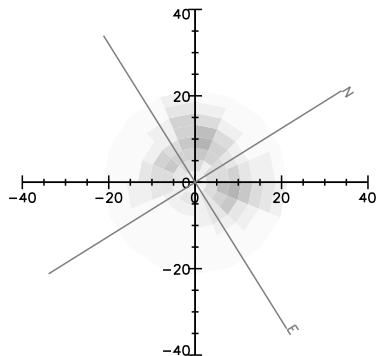
風速・風向の頻度分布



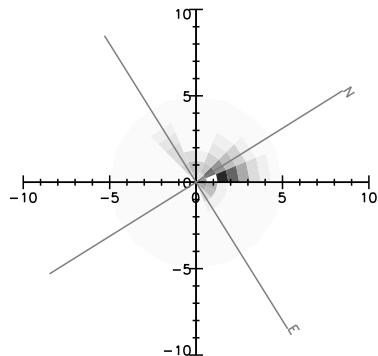
波高・波向の頻度分布



風速・風向の頻度分布

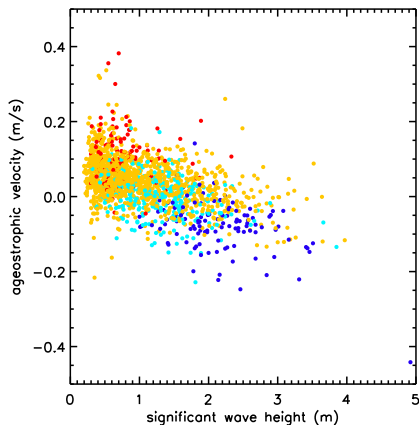


波高・波向の頻度分布

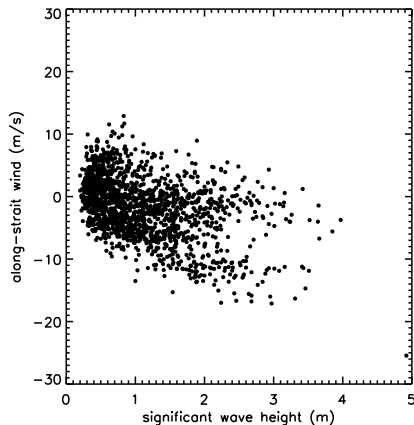


有義波高 vs 非地衡流、風速

有義波高 vs 非地衡流



有義波高 vs 風速



摩擦速度で規格化後、波浪依存性を調査する必要あり

非地衡流 と 吹送流 (x 成分) :

$$u_s^n = \text{表層流 (= 海洋レーダ一流速)}$$

$$u_g^n = \text{地衡流 (} \propto \text{水位差)}$$

$$= \frac{g}{fL}(\eta_1^n - \eta_2^n + \Delta\eta) = \frac{g}{fL}(\eta_1^n - \eta_2^n) + c$$

$$u_a^n = \text{非地衡流}$$

$$= u_s^n - u_g^n$$

$$u_w^n = \text{吹送流 (} \propto \text{摩擦速度または風速)}$$

$$= \alpha^m \cos \theta^m w_x^{n-h} + \alpha^m \sin \theta^m w_y^{n-h} = a^m w_x^{n-h} + b^m w_y^{n-h}$$

未知変数 :

$$\alpha^{m(n)} = m \text{ グループにおける風力係数}$$

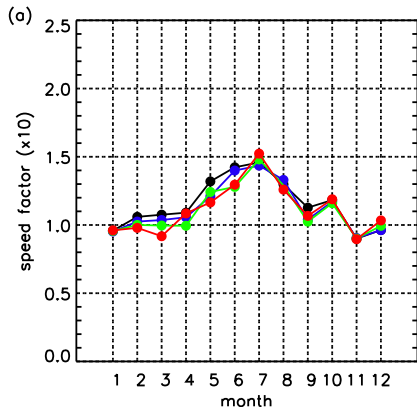
$$\theta^{m(n)} = m \text{ グループにおける吹送係数}$$

$$\Delta\eta = \text{基準点水位差}$$

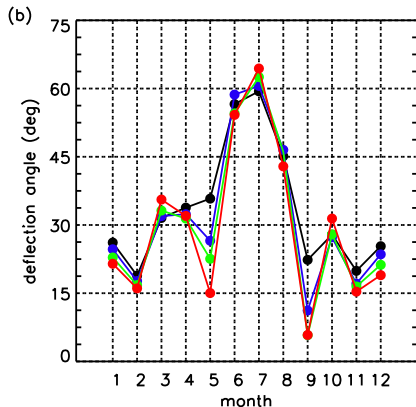
- 冬季（11月～2月）
成層の季節変化の影響を排除
- $|\theta^{wave} - \theta^{wind}| \leq 22.5 \text{ deg}$ （波向 $\simeq$ 風向）
波浪場の海峡代表性
- $H_S \geq 0.25 \text{ m}$
- $U_{10} \geq 2 \text{ m/s}$

結果：季節変化（摩擦速度基準）

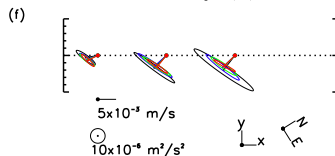
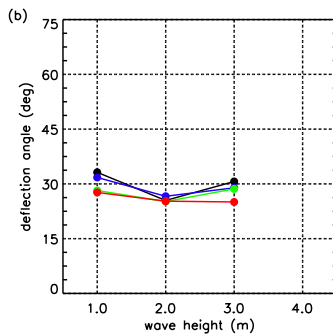
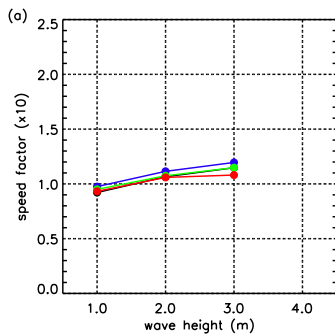
風力係数 : α^m



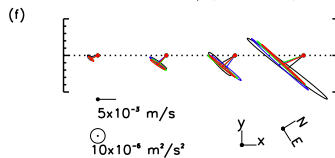
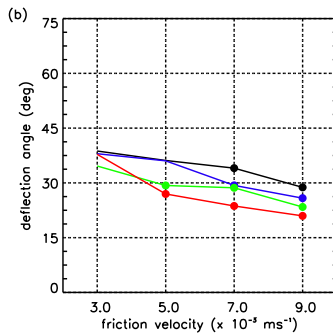
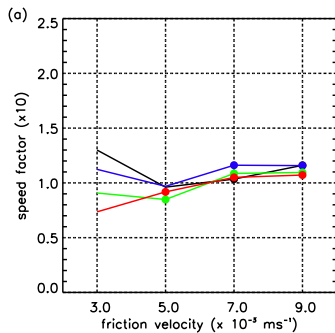
吹送偏角 : θ^m



風力係数と偏角：有義波高依存性

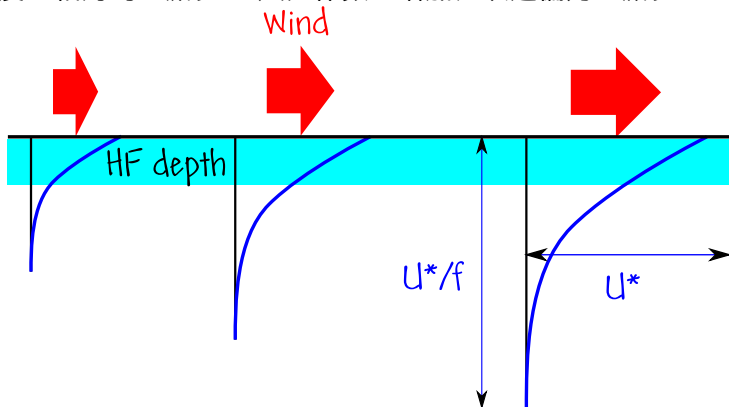


風力係数と偏角：摩擦速度（風速）依存性

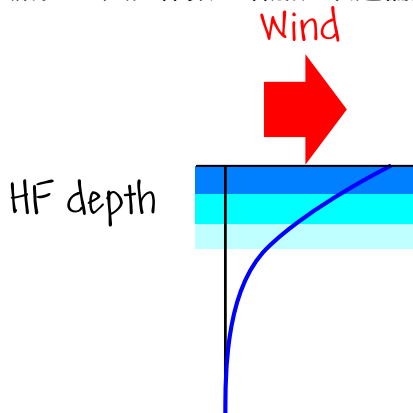


摩擦速度（風速）依存性

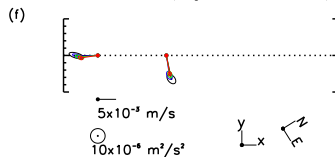
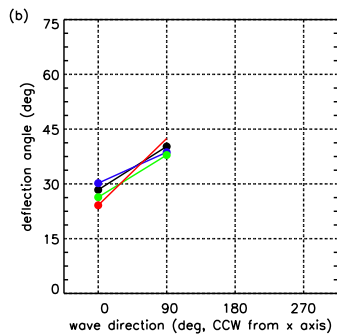
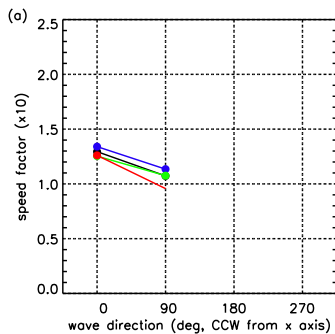
計測深度の相対的な減少 ⇨ 風力係数の増加、吹送偏角の減少



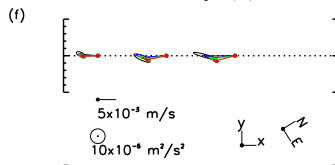
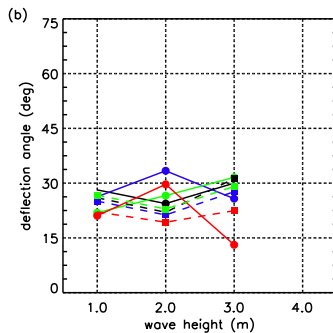
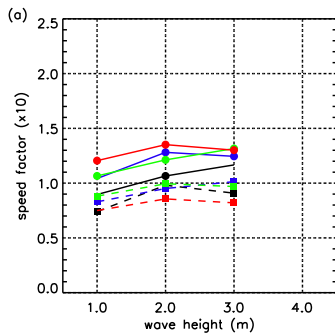
計測深度の相対的な減少 ⇨ 風力係数の増加、吹送偏角の減少



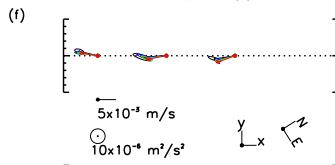
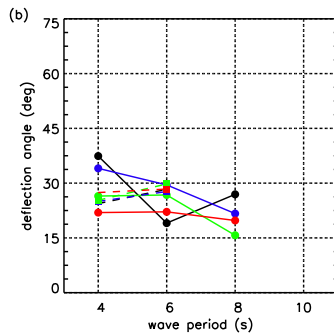
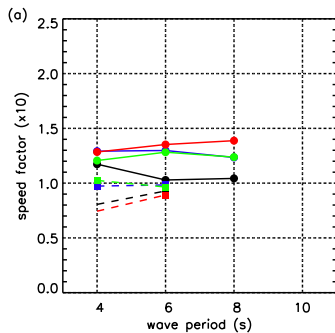
風力係数と偏角：波向依存性



風力係数と偏角：波向・有義波高依存性



風力係数と偏角：波向・有義周期依存性



① ストークスドリフト

- $u_{ST}^{R,V} \neq u_{ST}^{IR,IV}$
- $u_{ST}^{R,V}$ も U_{EK} も U_{10} に依存するので、切り分けが難しい。

② 対馬海峡の表層流の波浪依存性

- 波向依存性が見える
吹送距離の違いを反映？
- 有義波高依存性は見えるが、因果関係は不明
- 有義周期（波長）依存性は見えない

③ 波浪データの拡充を予定

④ 外洋も